

The impact of massive neutrinos on the large-scale structure of the Universe



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Neutrinos and Cosmology

Neutrino total mass from beta-decay experiments

$$0.05 {\rm eV} \lesssim \sum_{i} m_{\nu_i} \lesssim 2 \, {\rm eV}$$
 [T. Thummler et al. (2010)]
ArXiv:1012.2282

 Massive neutrinos have an impact on cosmological observables

 Cosmology can help in tightening the constraints on the neutrino total mass

Neutrinos and Cosmology

Neutrino total mas

 $0.05 {\rm eV} \lesssim$

 Massive neutrinos observables



 Cosmology can help in tightening the constraints on the neutrino total mass

Neutrinos and Cosmology

Neutrino total mass from beta-decay experiments

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Components of the Universe

Radiation

Components of the Universe

Ordinary Matter ('baryons')

Components of the Universe

The dark side:

Dark matter

Dark Energy (Λ)

ACDM model

- Expanding universe
- We describe Cold Dark Matter as an expanding pressureless fluid:



Otherwise pressure would prevent gravitational collapse!

(another reason for having CDM...)

Structures form! Stars, galaxies...

Isotropy and homogeneity

The APM Galaxy survey Maddox Sutherland Efstathiou & Loveday

Statistical properties

- Count number of objects (N) in spheres of radius R
- Define an overdensity field

 $\delta_R \equiv N/\bar{N} - 1$

- Autocorrelation: $\xi_R(r) = \langle \delta_R(\mathbf{x}) \delta_R(\mathbf{x} + \mathbf{r}) \rangle$
- Power Spectrum: $P(k_1)\delta_D(\mathbf{k}_1 + \mathbf{k}_2) = \langle \delta_{\mathbf{k}_1} \delta_{\mathbf{k}_2} \rangle$



The Power Spectrum



Massive neutrinos Effects on the clustering properties



Viel et al, JCAP (2010)

Massive neutrinos Effects on the power spectrum



Massive neutrinos Effects on the power spectrum



- Fermi-Dirac phase-space distribution → effective pressure
- Effective neutrino pressure **contrasts** the gravity-driven collapse at all scales smaller than a characteristic 'free streaming scale', $\lambda_{\rm fs}$, corresponding to a wavenumber $k_{\rm fs}$



 Damping of neutrino clustering at small scales → reflected also on CDM and galaxy clustering

Massive neutrinos Constraining their mass

• We could use the shape of the measured galaxy power spectrum but...

Galaxies are not all that 'matters'!!!

 Galaxy form only in the densest regions → they are a biased sampling of the total matter distribution

 $P_{\rm gal}(k) = F(P_{\rm tot\,matter}(k))$

Massive neutrinos Constraining their mass

- We can optimally choose to probe scales where non linearities (i.e. small scale physics) are not important
- The relation becomes a simple one!

 $P_{\text{gal}}(k) = b_{\text{lin}}^2 P_{\text{tot matter}}(k)$

 An observable designed specifically to exploit this, the clustering ratio, depends on the power spectrum and is unbiased at linear scales:

[Bel et al, A&A (2014)]

 $\eta_{\rm gal} \equiv \eta_{\rm tot\ matter}$

Want to know more about the clustering ratio? Just ask!



Search for the set of parameters of the model that maximises the likelihood

→ constraints on total neutrino mass



Search for the set of parameters of the model that maximises the likelihood

→ constraints on total neutrino mass



Search for the set of parameters of the model that maximises the likelihood

→ constraints on total neutrino mass





Conclusion

- Cosmology can help in tightening the constraints on neutrino mass
- The distribution of galaxies can be used for this purpose

Measuring its power spectrum (bias!!)

Using some smart observable (clustering ratio)

• What to do?

Theoretical characterisation of clustering ratio Measures of clustering ratio Monte-Carlo Markov Chain runs

. . .

Massive neutrinos Effects on the clustering properties

No neutrinos

Neutrinos

Matter clustering Fluid description

• Matter can be modelled as an expanding fluid, governed by:

$$\begin{cases} \frac{\partial \rho}{\partial t} + \nabla_{\mathbf{r}}(\rho \mathbf{u}) = 0\\ \frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla_{\mathbf{r}})\mathbf{u} = \frac{1}{\rho} \nabla_{\mathbf{r}} P + \nabla_{\mathbf{r}} \phi\\ \nabla_{\mathbf{r}}^2 \phi = 4\pi G \rho \end{cases}$$

 With (a bit more than) some algebra, these can be arranged into the linear equation of growth of fluctuations in an expanding fluid:

$$\ddot{\delta} + 2H\dot{\delta} = 4\pi G\bar{\rho}\delta + \frac{c_s^2}{a^2}\nabla^2\delta$$

Matter clustering Fluid description

 Cold Dark Matter is supposed to be pressureless, so we can safely neglect the speed of sound here...

$$\ddot{\delta} + 2H\dot{\delta} = 4\pi G\bar{\rho}\delta + \frac{c_s^2}{a^2}z_{\delta}^2$$

• ...but is it the same for neutrinos?

[spoiler: no!]

They are characterised by a Fermi-Dirac phase-space distribution:

$$f(\mathbf{x}, \mathbf{p}, t) = \frac{1}{e^{-\frac{\mathbf{p}}{T_{\nu}}} + 1}$$

Their density is

$$\rho(\mathbf{x},t) = \frac{g}{(2\pi)^3} \int d^3 \mathbf{p} f(\mathbf{x},\mathbf{p},t) E(\mathbf{p})$$

$$E = \sqrt{|\mathbf{p}|^2 + m_\nu^2}$$

And their pressure

$$P(\mathbf{x},t) = \frac{g}{(2\pi)^3} \int d^3 \mathbf{p} f(\mathbf{x},\mathbf{p},t) \frac{|\mathbf{p}|^2}{E(\mathbf{p})}$$

From density and pressure one can find their effective sound speed

$$c_s = 134.423(1+z) \left[\frac{1 \,\mathrm{eV}}{m_{\nu}}\right] \,\mathrm{km/s}$$

 This is not negligible! When considering massive neutrinos, we will have to solve a system of equations, one for CDM and one for neutrinos → 2 fluid approximation

• 2 fluid approximation:

$$\begin{cases} \ddot{\delta}_{\rm c} + 2H\dot{\delta}_{\rm c} - \frac{3}{2}H^2\Omega_m \left\{ [1-\nu]\delta_{\rm c} + \nu\delta_\nu \right\} \\ \ddot{\delta}_\nu + 2H\dot{\delta}_\nu - \frac{3}{2}H^2\Omega_m \left\{ [1-\nu]\delta_{\rm c} + [\nu - (k/k_{\rm fs})^2]\delta_\nu \right\} \end{cases}$$



Statistical properties

- Count number of objects (N) in spheres of radius R
- Define an overdensity field

 $\delta_R \equiv N/\bar{N} - 1$

- Spatial average $\langle \delta_R \rangle = 0$
- Variance $\sigma_R^2 = \langle \delta_R^2 \rangle$



• Autocorrelation $\xi_R(r) = \langle \delta_R(\mathbf{x}) \delta_R(\mathbf{x} + \mathbf{r}) \rangle$

The Clustering Ratio Motivation



The Clustering Ratio Motivation

But they are not scorrelated!

$$\delta_{g,R}(\mathbf{x}) = F[\delta_{m,R}(\mathbf{x})]$$

 Smoothing on linear scale (and assuming this function to be local)

$$\delta_{g,R}(\mathbf{x}) = b_{lin} \delta_{m,R}(\mathbf{x})$$
$$\xi_{g,R}(r) = b_{lin}^2 \xi_{m,R}(r)$$
$$\sigma_{g,R}^2 = b_{lin}^2 \sigma_{m,R}^2$$

The Clustering Ratio Definition

• Therefore, if we define the clustering ratio as...

$$\eta = rac{\xi(r)}{\sigma^2}$$
 [Bel et al.(2014)]

• ... it ends up having this very good property:



The Simulations DEMNUni

- Dark Energy and Massive Neutrino Universe
 [C. Carbone et al. (2015) in prep, E. Castorina et al, JCAP (2015)]
- Set of 4 simulations with Planck13 cosmology and neutrino mass = { 0, 0.17, 0.30, 0.53 } eV
- CDM mass resolution: ~8x10¹⁰
 Number of CDM particles: 2048³
 Number of neutrino particles: 2048³
 Cubical box of side: 2000 h

~8x10¹⁰ *h* ⁻¹ M_{sun} 2048³ 2048³ 2000 *h* ⁻¹ Mpc

• Run at CINECA by C. Carbone (5x10⁶ cpu hours) using the gadget - III code [Springel et al. (2005), Viel et al. (2010)]